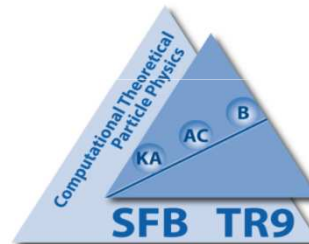


Higgs production via vector-boson fusion at NNLO in QCD

Paolo Bolzoni



In collaboration with F. Maltoni, S. Moch, M. Zaro [arXiv:1003.4451\[hep-ph\]](https://arxiv.org/abs/1003.4451)

Loopfest IX , 21-23 June 2010, Stony Brook University

Introduction and motivations

Motivations for NNLO

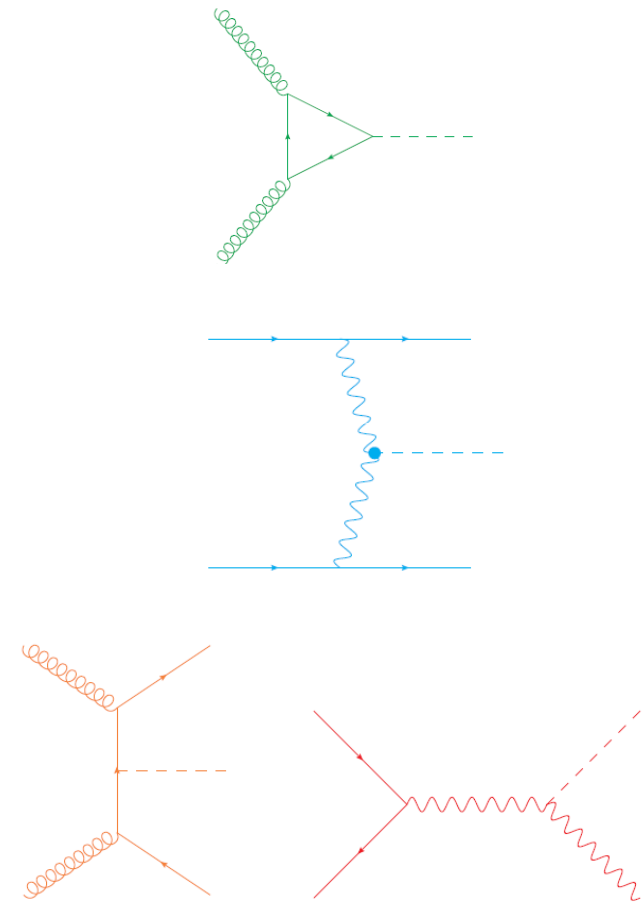
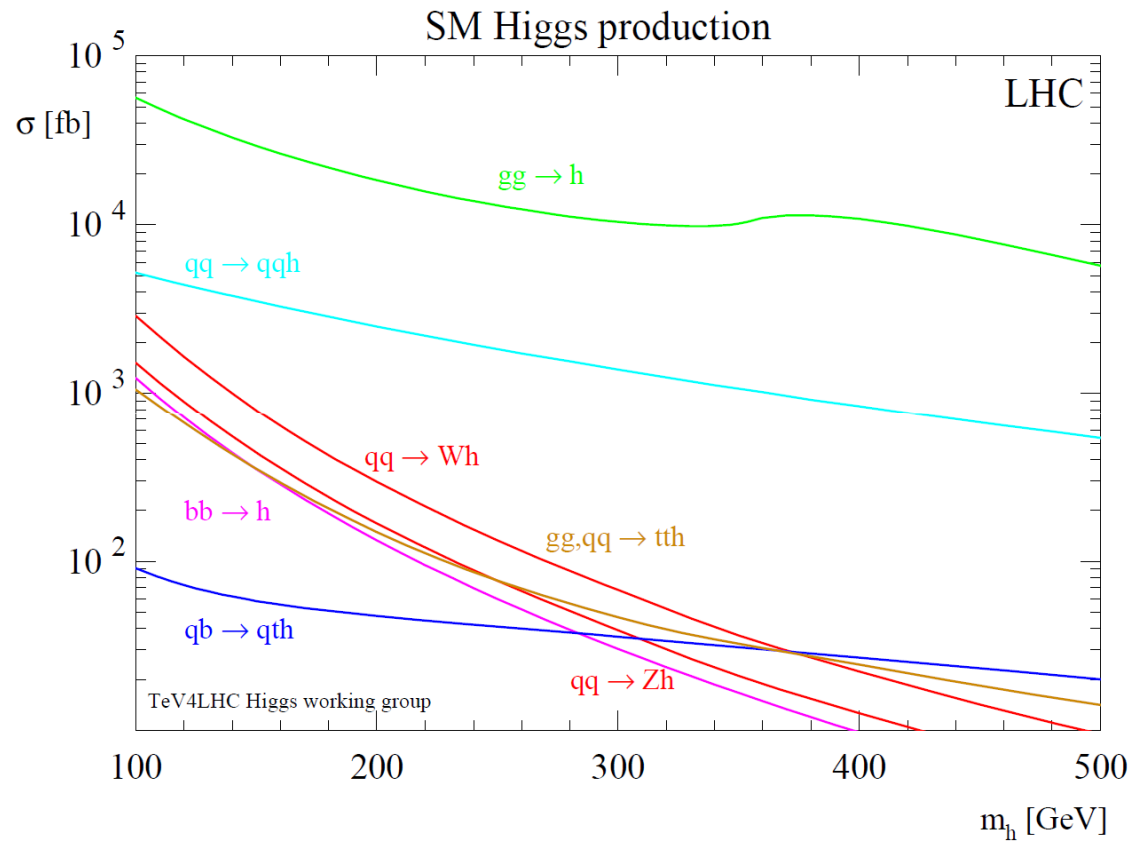
The search for the Higgs demands:

- **Its existence**
- **An estimate of the expected events**
- **Control over the background noise**

Then to investigate its properties we need:

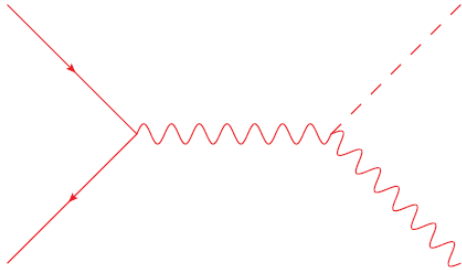
- **A determination of the cross section as precise as possible**
- **The NLO already gives a reliable estimation of the cross section**
- **As will be shown at NNLO there is a significant improvement on the theoretical uncertainties from 10-15% down to 1-2%**

Higgs boson production channels



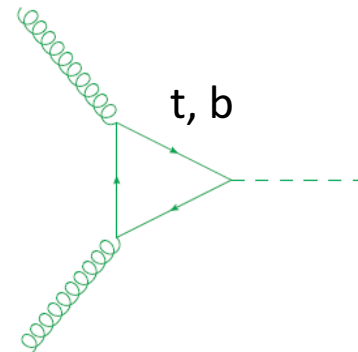
VBF is the second most important production mechanism

Higgs-Strahlung and gluon fusion at NNLO



Higgs-Strahlung computed at NNLO

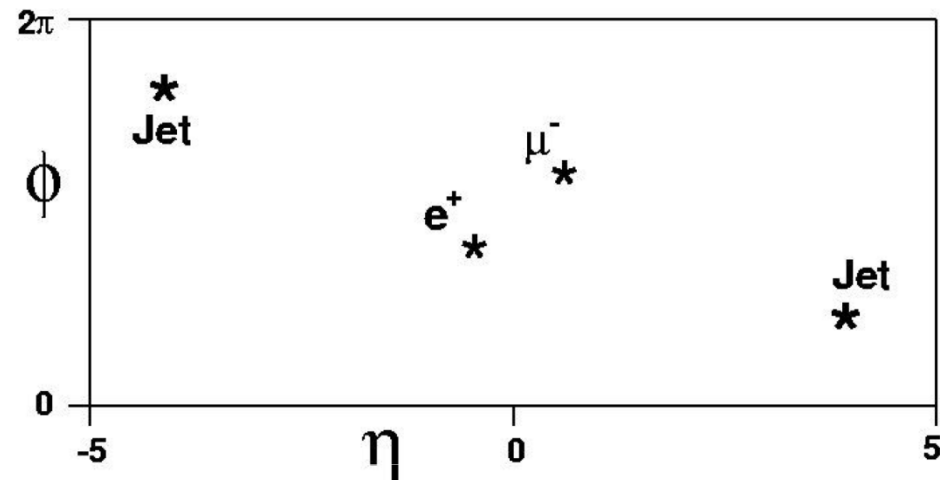
[O. Brein, A. Djouadi, R.V. Harlander (2004)]



Latest updates of gluon fusion at NNLO with the inclusion of finite top mass effect.

[A. Pak, M. Rogal, M. Steinhauser; R.V. Harlander, H. Mantler, S. Marzani, K.J. Ozeen]

The VBF signal

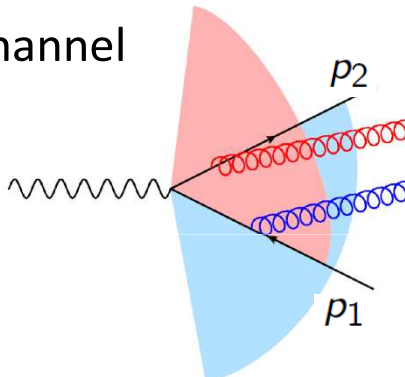


- Two hard tagging jets
- Large rapidity separation between jets
- Very small hadronic activity between jets
- Higgs decay in the central rapidity region

Hadronic activity in VBF

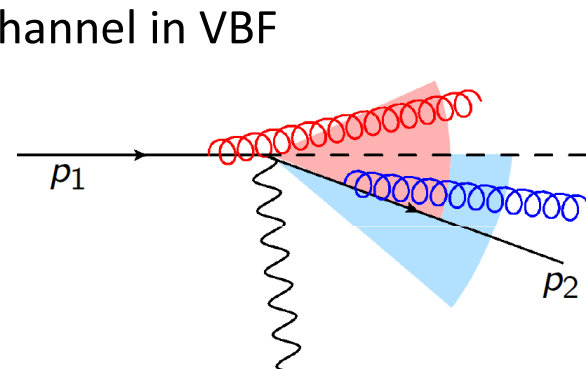
Angular ordering for the dominant soft gluon emission: the averaged azimuthal emission is confined in a cone given by the angle between the emitter and the spectator

s-channel



$$\sim \frac{1}{(p_1 + p_2)^2} = \frac{1}{s} = \frac{1}{4E_{cm}^2}$$

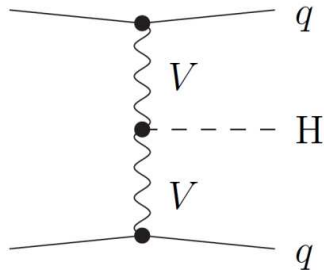
t-channel in VBF



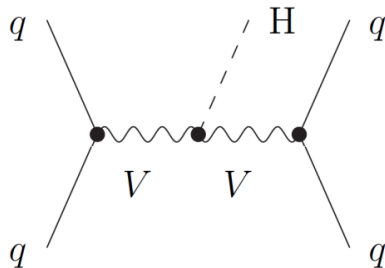
$$\sim \frac{1}{(p_1 - p_2)^2 - M_V^2} = \frac{-1}{M_V^2 + 2E_1 E_2 (1 - \cos \vartheta_{12})}$$

- The t-channel propagator dominates for small ϑ_{12}
- Soft gluon radiated at $\vartheta_g < 2\vartheta_{12}$

The VBF is a well defined process by itself

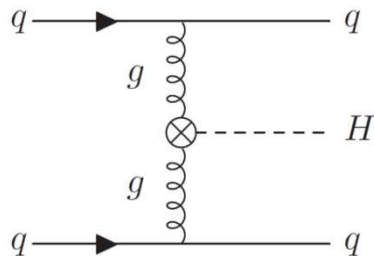


The VBF has interferences with other processes at LO but also at higher orders



The interference with the Higgs associated production is at the **per mil level already at LO in QCD**

[M. Ciccolini, A. Denner, S. Dittmaier]



At higher orders the interference with the gluon fusion process is **well below the percent level**

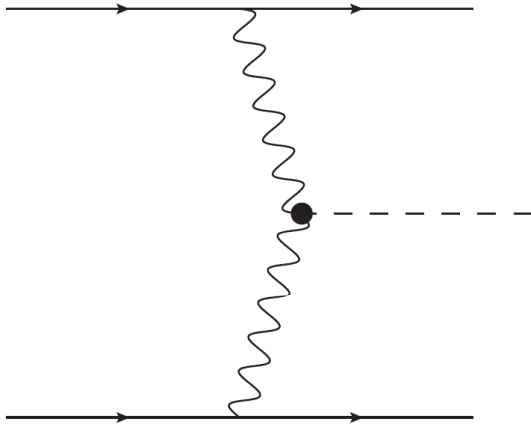
[J.R. Andersen, T.Binoth, G. Heinrich, J.M. Smillie]

VBF can be defined up to an ambiguity of 1% which sets the target precision for theoretical predictions

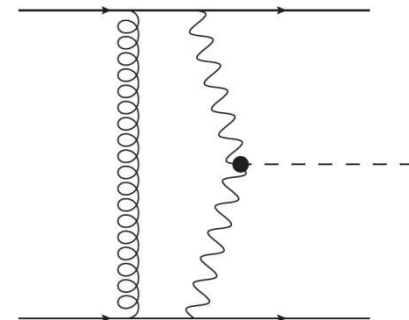
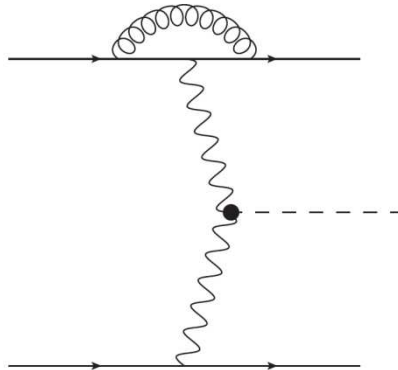
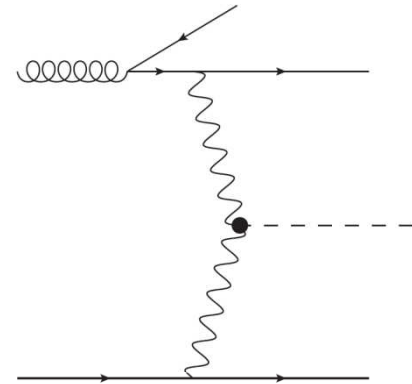
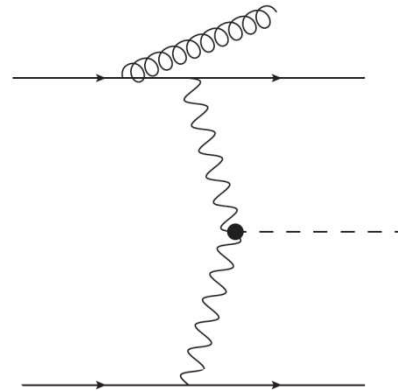
The computation

The structure function approach at NLO

LO

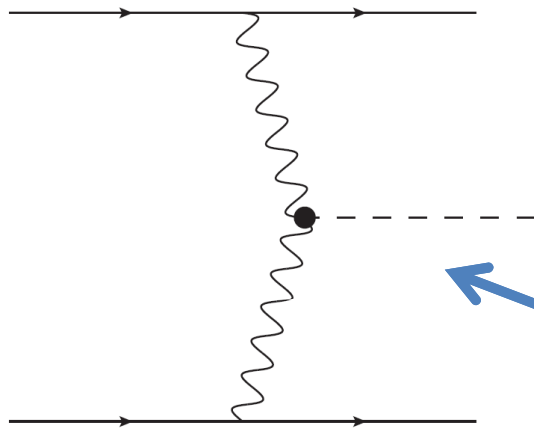


NLO corrections

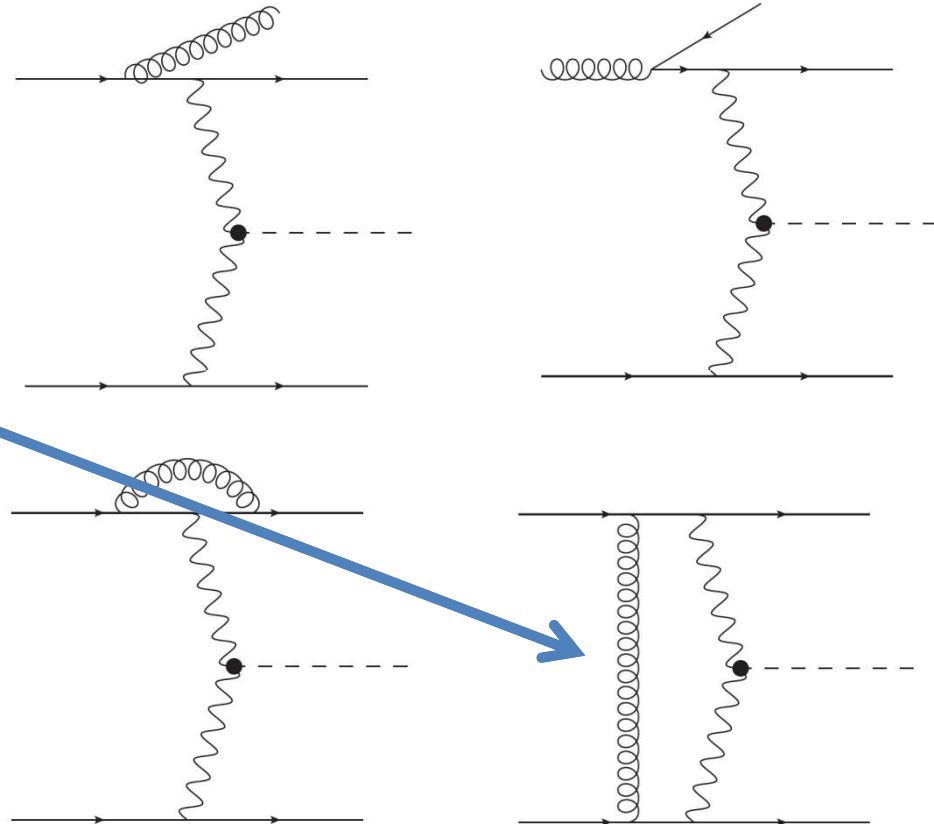


The structure function approach at NLO

LO



NLO corrections

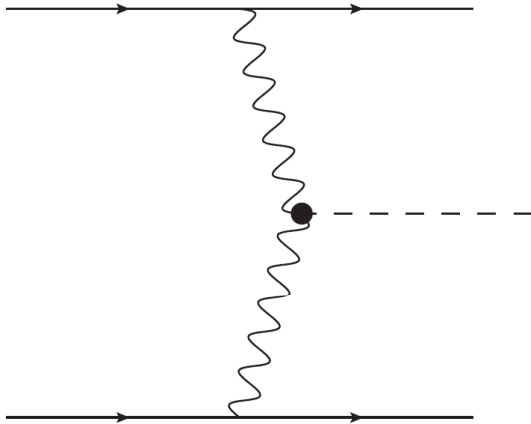


The interference of the last type of diagrams with the LO has a vanishing color factor:

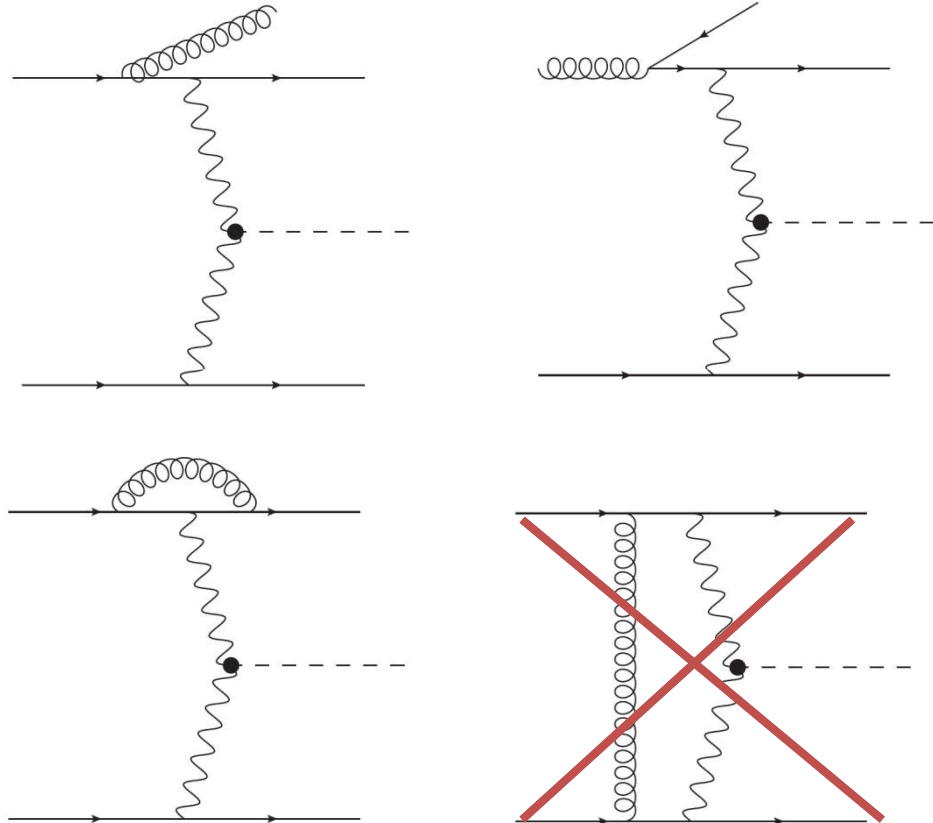
$$\delta_{ij} t_{ji}^a \delta_{ml} t_{lm}^a = [\text{Tr}(t)]^2 = 0$$

The structure function approach at NLO

LO

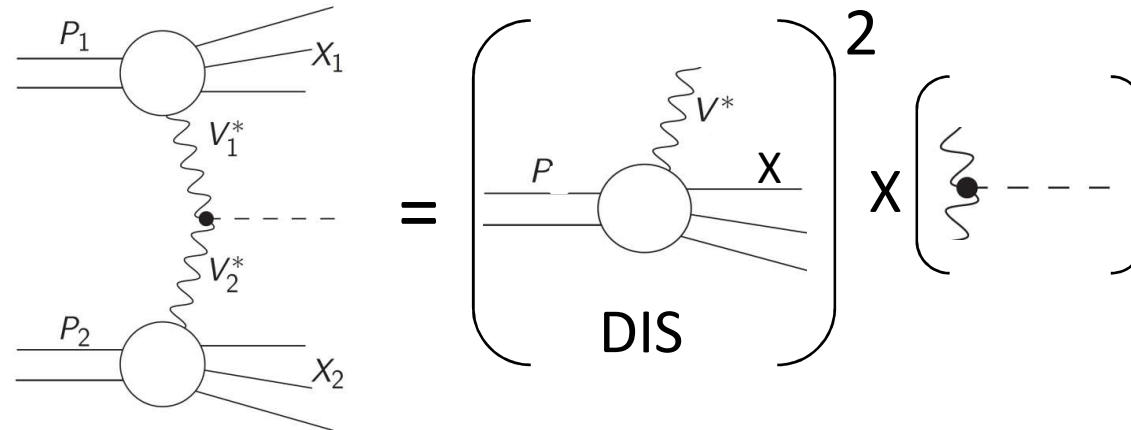


NLO corrections



Hence we can neglect the diagrams with a gluon exchange in the t-channel and the upper and lower lines remain independent

The structure function approach at NLO cntd.



$$d\sigma = \frac{1}{S} G_F^2 M_{V_1}^2 M_{V_2}^2 \frac{1}{(Q_1^2 + M_{V_1}^2)^2} \frac{1}{(Q_2^2 + M_{V_2}^2)^2} W_{\mu\nu}^{DIS}(x_1, Q_1^2) \mathcal{M}_{V VH}^{\mu\rho} \mathcal{M}_{V VH}^{*\nu\sigma} W_{\rho\sigma}^{DIS}(x_2, Q_2^2) \\ \times \frac{d^3 P_{X_1}}{(2\pi)^3 2E_{X_1}} \frac{d^3 P_{X_2}}{(2\pi)^3 2E_{X_2}} ds_1 ds_2 \frac{d^3 P_H}{(2\pi)^3 2E_H} (2\pi)^4 \delta^4(P_1 + P_2 - P_{X_1} - P_{X_2} - P_H)$$

Hence the NLO knowledge of the DIS structure functions F_1 , F_2 and F_3 is enough to compute the NLO VBF cross section

This factorization exact at NLO is the so-called **structure function approach**

[T.Hahn, G. Valencia, S. Willenbrock]

The structure function approach at NNLO

At NNLO the structure function approach is NOT exact but it can **be still considered a good approximation** and the dominant contributions can be included in the NNLO structure functions

$$F_1, F_2, F_3$$

[D.I.Kazakov, A.V. Kotikov, G.Parente, O.A.Sampaio, J.S.Guillen,
E.B.Zijlstra, W.L. van Neerven, S.Moch, J.A.M.Vermaseren,A.Vogt]

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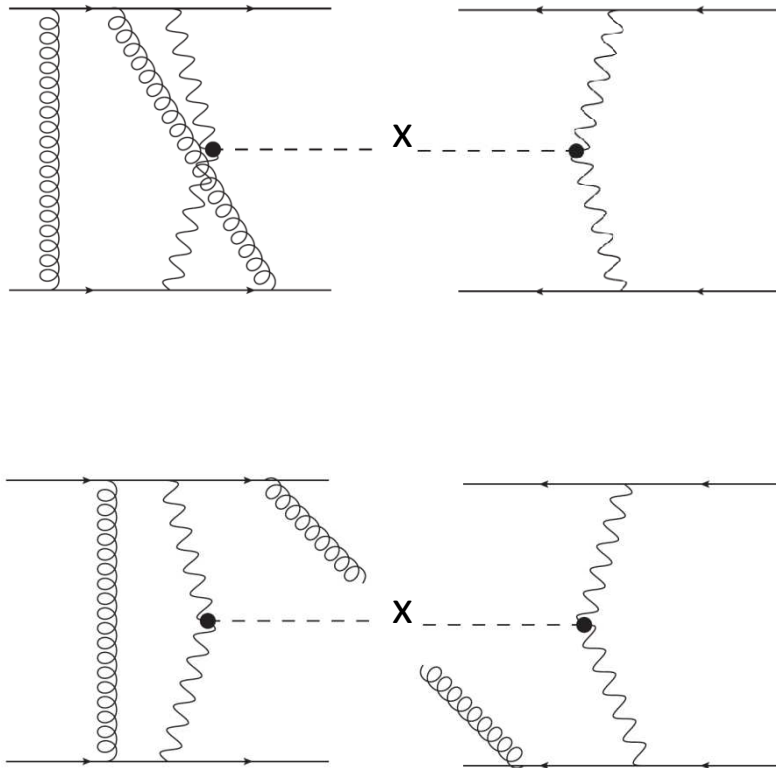
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[D.I.Kazakov, A.V. Kotikov, G.Parente, O.A.Sampaio, J.S.Guillen,
E.B.Zijlstra, W.L. van Neerven, S.Moch, J.A.M.Vermaseren,A.Vogt]

The following types of contributions that in principle destroy the structure function approach are:

- Double gluon-exchange in the t-channel
(real and virtual)
- Single quark line (SQL) diagrams
- t/b loop diagrams

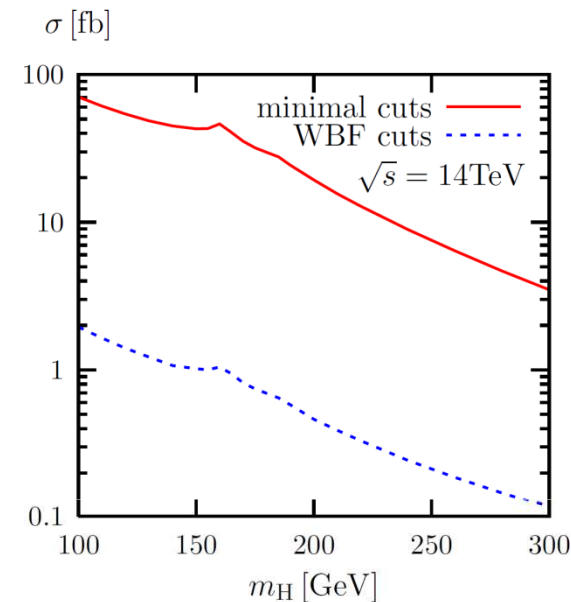
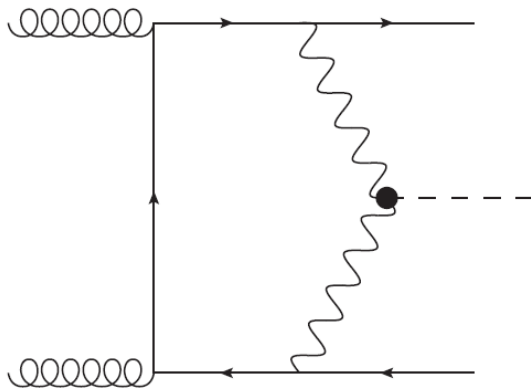
The double-gluon exchange diagrams



- This is a gauge invariant class of diagrams
- IR and UV finite
- It is a color $(1/N_c)^2$ suppressed class with respect to the DIS^2 contributions
- They are also kinematically suppressed

[T.Figy, V. Hankele, D.Zeppenfeld,
J. Blümlein, J.A.M. Vermaseren]

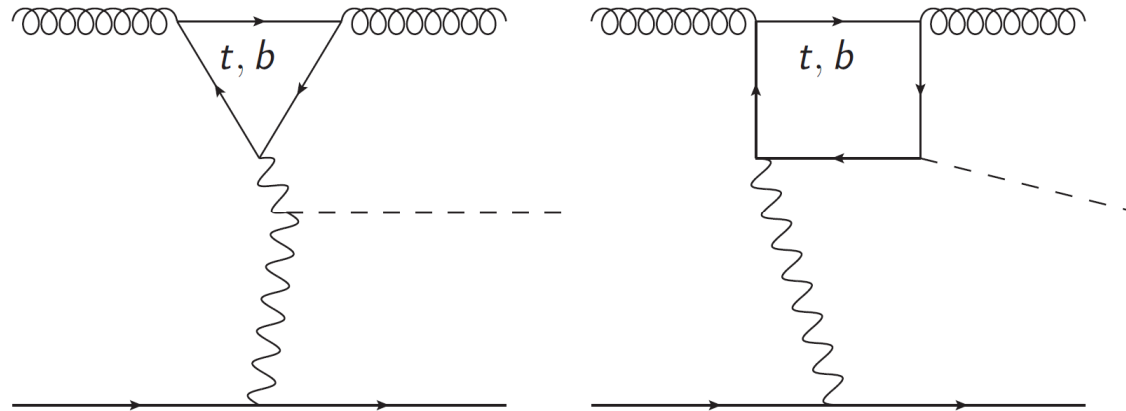
The SQL diagrams



Single quark line (SQL) contributions: “using a minimal set of cuts, the numerical impact of these terms is at the percent level with respect to NLO rate for WBF. Applying the so-called WBF cuts leads to an even stronger suppression, so that **we do not expect a significant deterioration of the WBF** signal by these color exchange effects”

[R.V. Harlander, J. Vollinga, M.M. Weber arXiv:0801.3355[hep-ph]]

The t/b loop diagrams

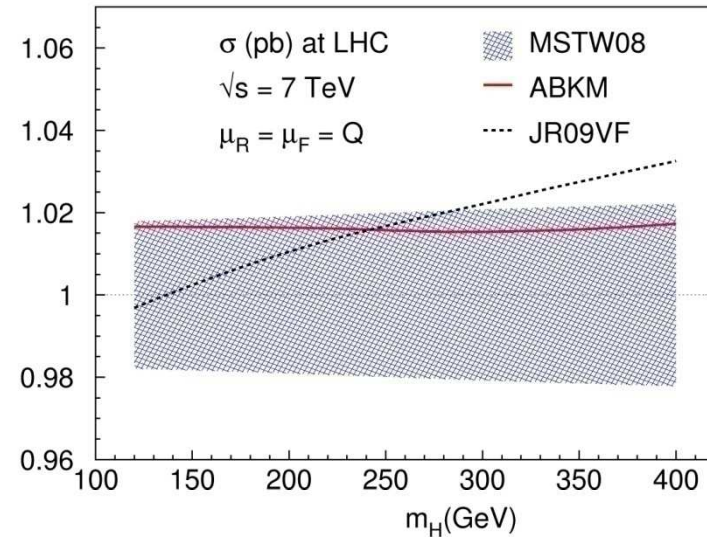
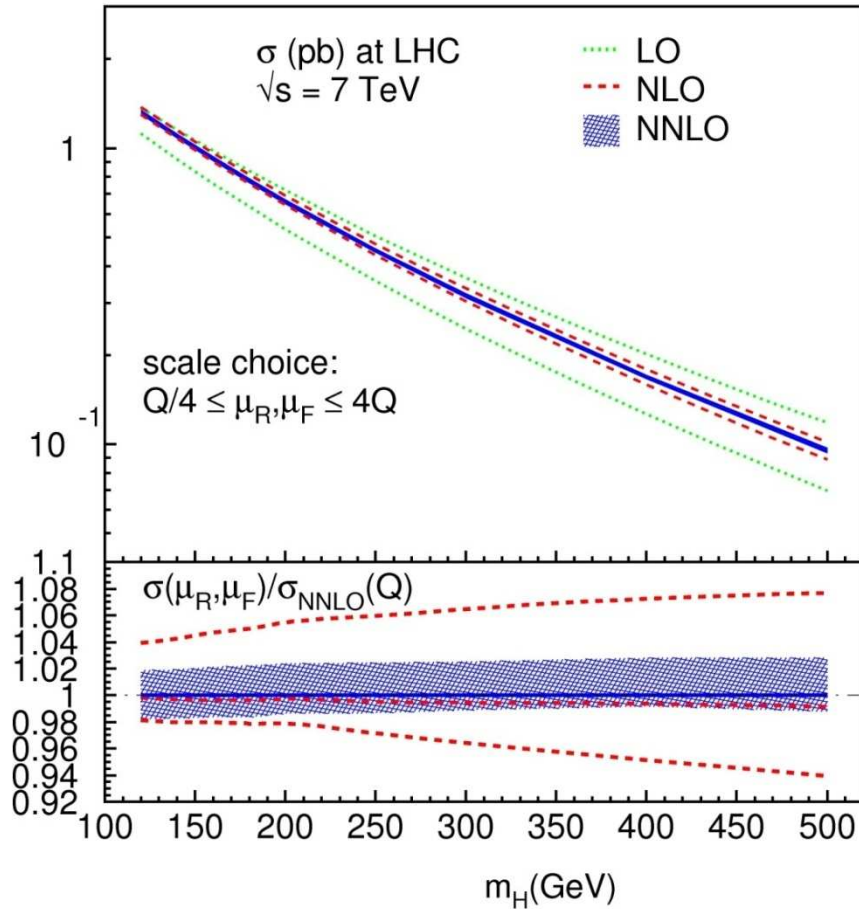


m_h (GeV)	120	300	500
1.96 TeV	3.87 E-6 (0.0690)	2.52 E-7 (0.0054)	1.50 E-8 (0.00042)
7 TeV	2.62 E-4 (1.235)	7.89 E-5 (0.614)	2.73 E-5 (0.088)

- **Contribution from the triangle to the total VBF cross section well below the 1% (checked by two independent computations)**

Results

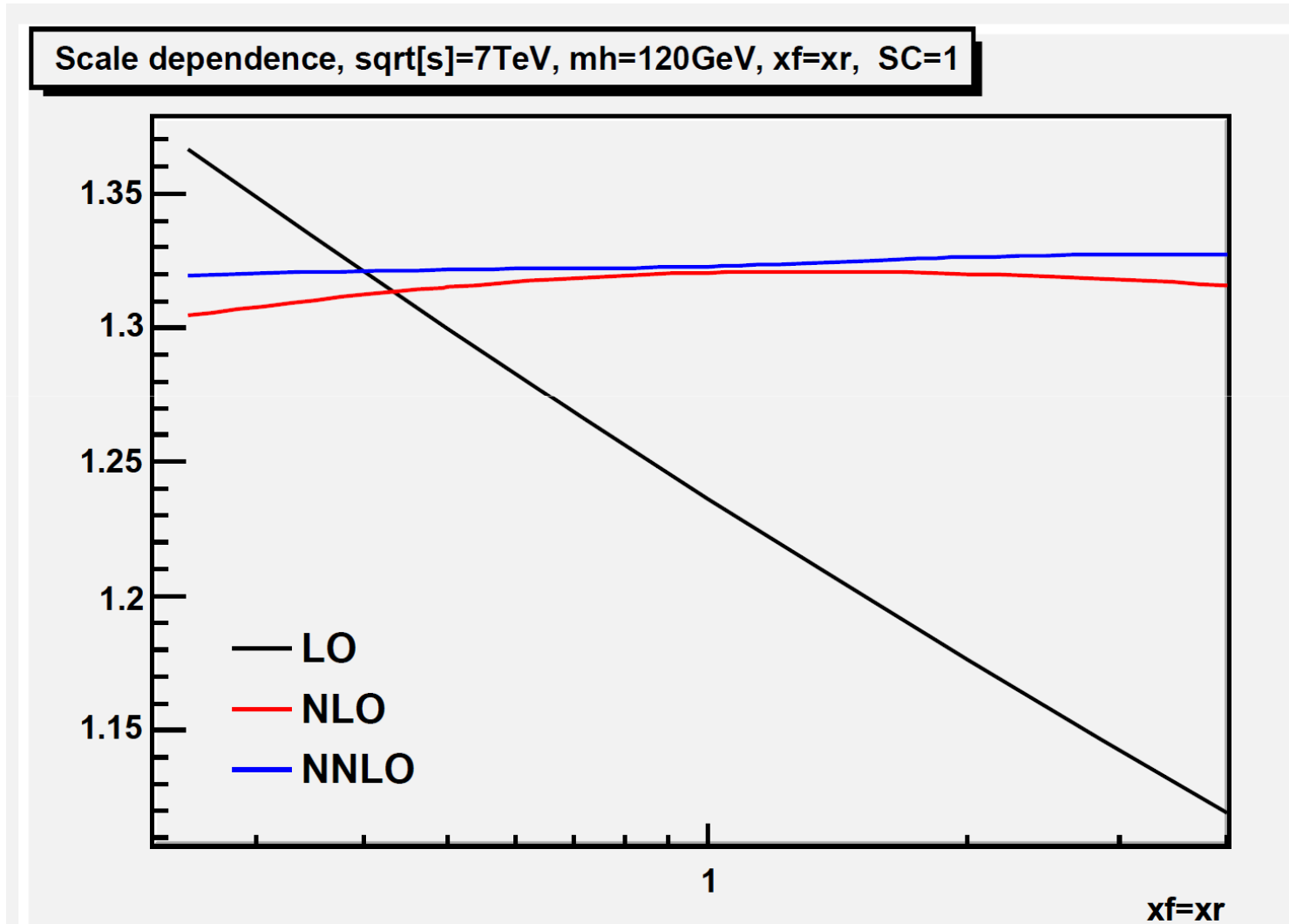
LHC @ 7 TeV



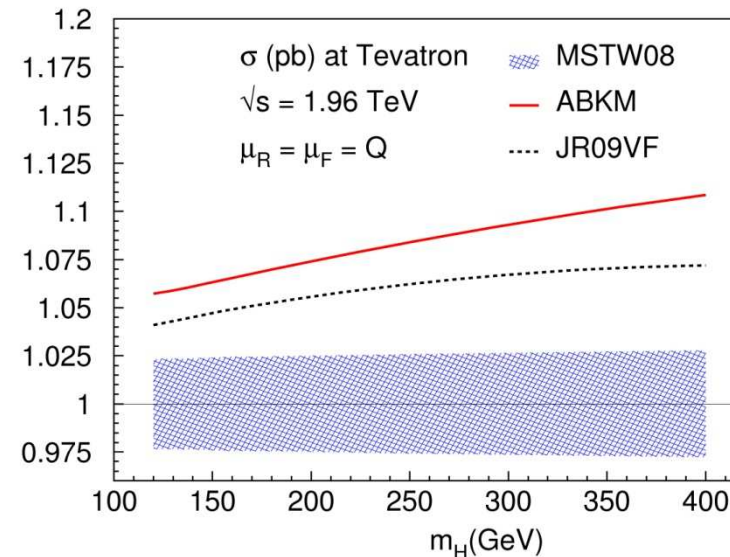
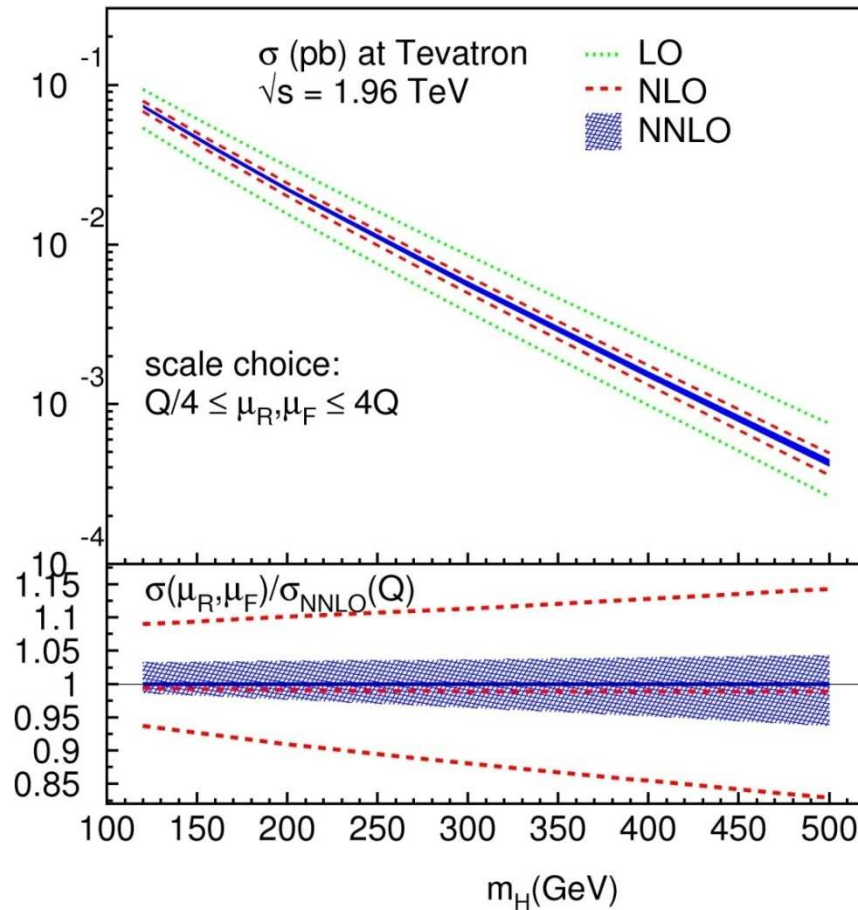
The PDF error band MSTW08 is 2% and is compatible with other NNLO best fits (ABKM, JR)

The convergence of the perturbative series is well behaved

Theoretical uncertainties improvement



Tevatron



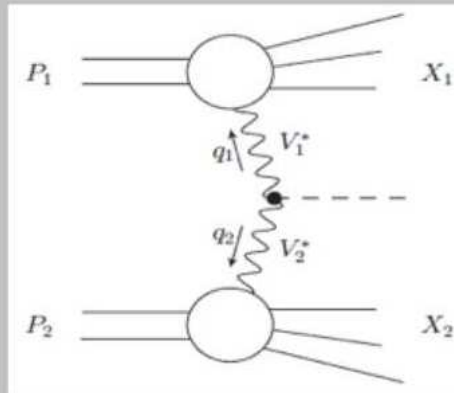
In this case the difference between MSTW08 and ABKM and JR is due to larger uncertainties for the high- x quark PDFs

[S. Alekhin, J. Blümlein, S. Klein, S. Moch arXiv:0908.2766[hep-ph]]

VBF @ NNLO : Cross-section Calculator

by P. Bolzoni, F. Maltoni, S.-O. Moch and M. Zaro

alpha version v0.1 -- 10 April 2010



Higgs production in vector-boson fusion (VBF) is computed via a structure-function approach, as reported in [ArXiv:1003.4451 \[hep-ph\]](https://arxiv.org/abs/1003.4451).

This simple interface allows any [registered](#) user to obtain a cross section up to NNLO in QCD, including an estimate of the theoretical uncertainties coming from scale variation and PDF uncertainties.

The electro-weak parameters used for the cross-section computation are set to their respective [PDG](#) values (see the list [here](#)).

The code runs over the CP3-MadGraph cluster and might take up to a few hours depending on the actual request. An e-mail with the corresponding data file is sent to the user as soon as results are available. The possibility of requesting multiruns, i.e. runs corresponding to a series of Higgs mass values and/or collider energies, will be available soon upon e-mail request.

See the [HNNLO web page](#) by M. Grazzini for a similar tool for $gg \rightarrow H$.

New users are kindly asked to register [here](#)

Up to order:	<input type="button" value="NNLO"/>	
Collider type:	<input type="button" value="p-p"/>	
Center of mass energy:	<input type="text" value="7000"/>	GeV
Higgs boson mass:	<input type="text" value="120"/>	GeV
PDF set:	<input type="button" value="MSTW08"/>	
PDF uncertainties:	<input type="button" value="no"/>	
Reference scale:	<input type="button" value="Q"/>	Description
Scale uncertainties:	<input type="button" value="no"/>	Description

<http://madgraph.phys.ucl.ac.be/vbf.html>

Conclusions and outlooks

Conclusions and outlooks

- **VBF is a promising channel for both the discovery and the precision measurement of the Higgs at LHC**
- **The structure function approach is a good tool to obtain NNLO cross section**
- **First NNLO VBF cross section**
- **Theoretical uncertainties lowered to 1-2%**
- **Need for a differential NNLO computation for the future**

Extra slides

Theoretical uncertainties estimate

DIS structure functions:

$$F_i(x, Q^2) = \sum_{j=0} \alpha_s^j(\mu_r^2) [C_i^j(Q^2, \mu_r^2, \mu_f^2) \otimes f(\mu_f^2)](x)$$

- Q^2 is the physical scale of the process
(the DIS cross section depends on it)
- $\mu_{f,r}$ unphysical factorization and renormalization scales

The theoretical uncertainties are estimated varying the two unphysical scales. This dependence tends to be compensated and hence the dep. of a order α^k computation is of order α^{k+1}

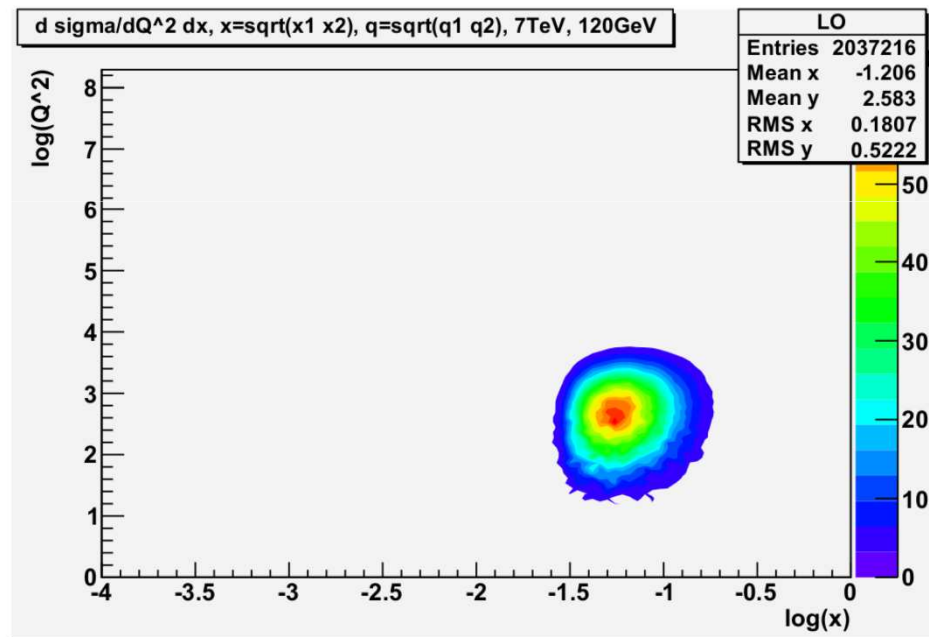
Scale choices

Two different scales have been used:

$$\mu_{r/f}^2 = x_{r/f} m_h^2 \quad (\text{SC} = 0)$$

$$\mu_{r/f}^2 = x_{r/f} Q^2 \quad (\text{SC} = 1)$$

$$\frac{1}{4} \leq x_{r/f} \leq 4$$

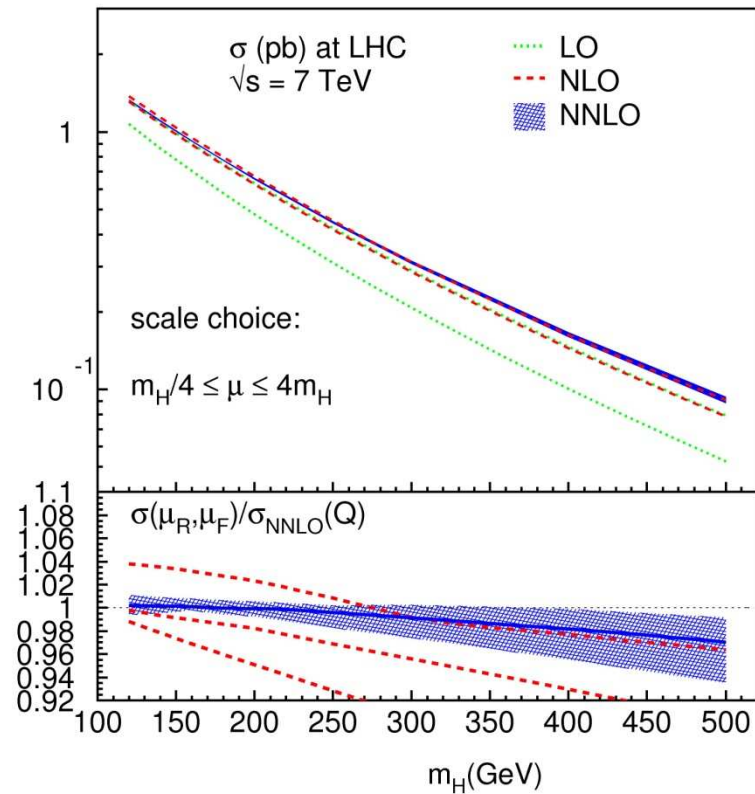
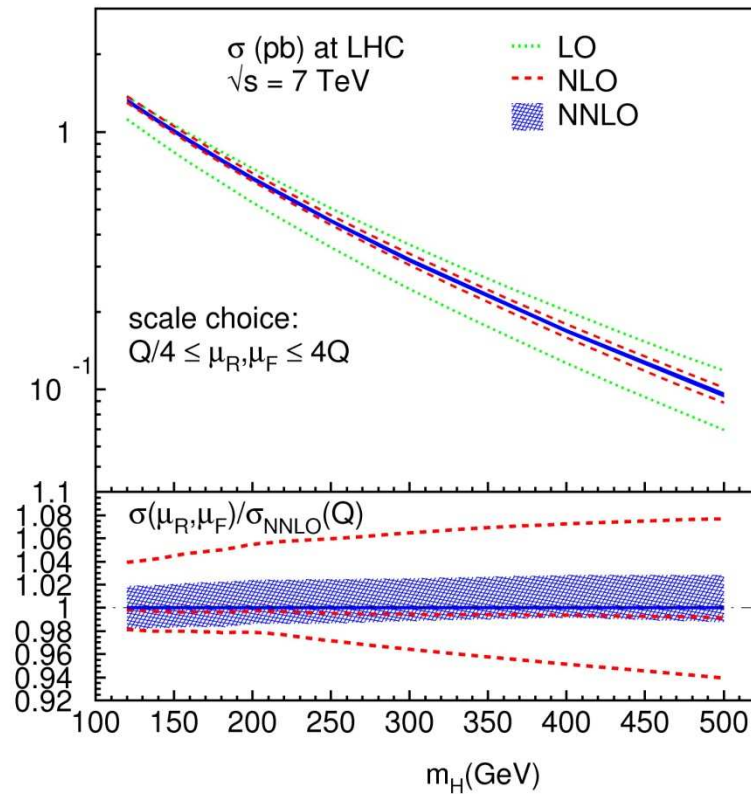


$$\langle Q^2 \rangle \simeq (20\text{GeV})^2$$

$$\langle x \rangle \simeq 6 \cdot 10^{-2}$$

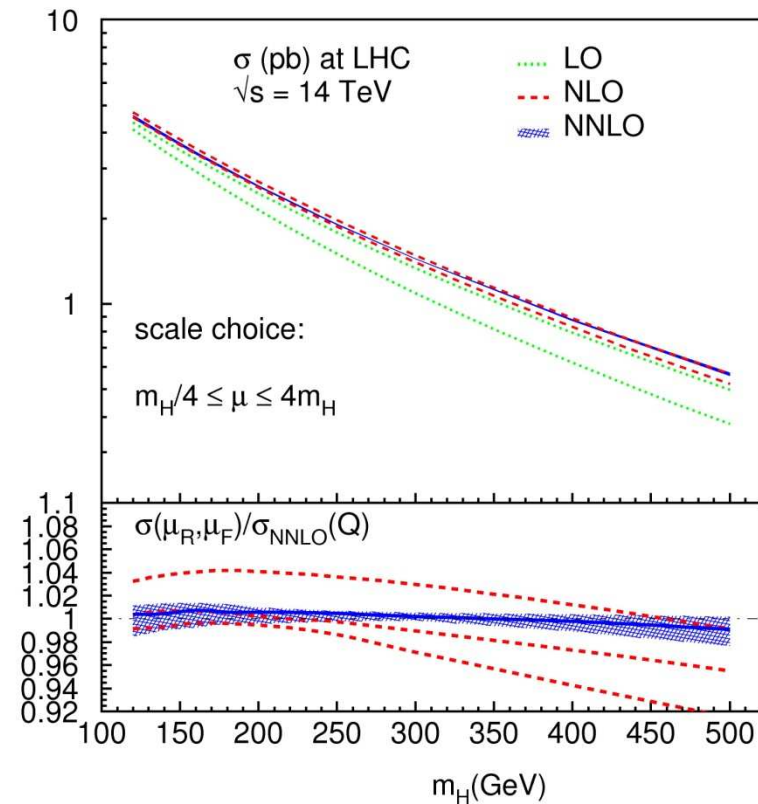
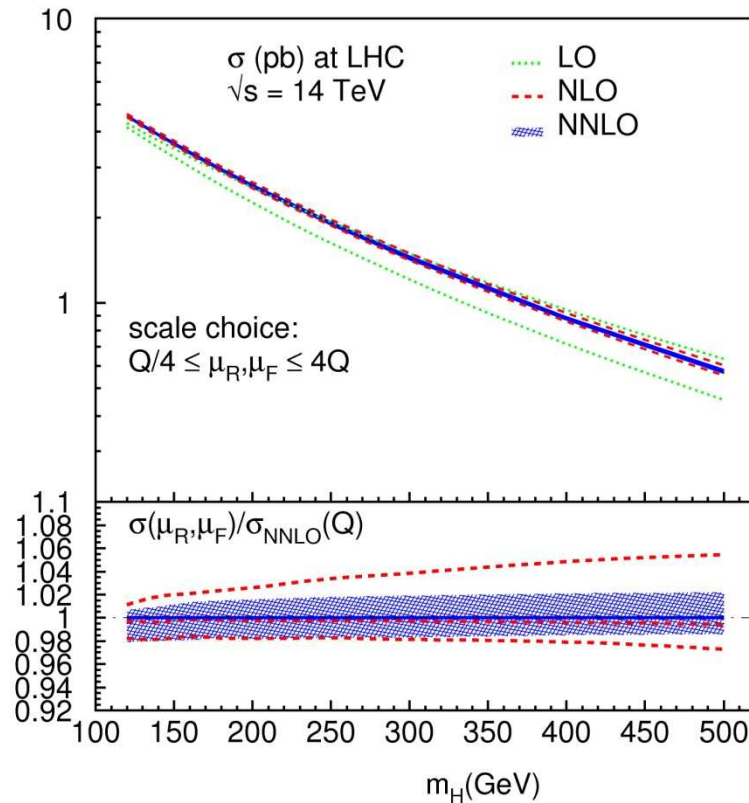
The effective averaged value of Q^2 is relatively low and so the choice $\text{sc}=1$ looks more natural

At NNLO the two scale choices become equivalent



LHC @ 7 TeV

At NNLO the two scale choices become equivalent



LHC @ 14 TeV

PDF uncertainties LHC @ 14 TeV

